

Derivation of the time dependent Hartree (Fock) equation

Peter Pickl

Mathematical Institute
LMU

13. April 2016

Dictionary

1. “Derivation” means prove of validity
2. “Hartree (Fock)”: time dependent Hartree (Fock)

Dictionary

1. “Derivation” means prove of validity
2. “Hartree (Fock)”: time dependent Hartree (Fock)

Dictionary

1. “Derivation” means prove of validity
2. “Hartree (Fock)”: time dependent Hartree (Fock)

Overview

1. Derivation of Hartree equations for Bosons
2. Derivation of Hartree (Fock) equations for Fermions
3. Special case: A tracer particle in the fermi sea

Overview

1. Derivation of Hartree equations for Bosons
2. Derivation of Hartree (Fock) equations for Fermions
3. Special case: A tracer particle in the fermi sea

Overview

1. Derivation of Hartree equations for Bosons
2. Derivation of Hartree (Fock) equations for Fermions
3. Special case: A tracer particle in the fermi sea

Overview

1. Derivation of Hartree equations for Bosons
2. Derivation of Hartree (Fock) equations for Fermions
3. Special case: A tracer particle in the fermi sea

Mean field for the bosons: The Hartree equation

$$H = \sum_{j=1}^N -\Delta_j + \sum_{j=1}^N A_t(x_j) + (N-1)^{-1} \sum_{k < j} V(x_j - x_k)$$

Interaction „felt“ by each particle of order one

$$\Psi_0 = \prod_{j=1}^N \phi_0(x_j) \quad id_t \Psi_t = H_t \Psi_t$$

Interaction destroys product structure.

Question:

- ▶ In which regimes: $\Psi_t \approx \prod_{j=1}^N \phi_t(x_j)$ (in what sense?)
- ▶ What is ϕ_t ?

Mean field for the bosons: The Hartree equation

$$H = \sum_{j=1}^N -\Delta_j + \sum_{j=1}^N A_t(x_j) + (N-1)^{-1} \sum_{k < j} V(x_j - x_k)$$

Interaction „felt“ by each particle of order one

$$\Psi_0 = \prod_{j=1}^N \phi_0(x_j) \quad id_t \Psi_t = H_t \Psi_t$$

Interaction destroys product structure.

Question:

- ▶ In which regimes: $\Psi_t \approx \prod_{j=1}^N \phi_t(x_j)$ (in what sense?)
- ▶ What is ϕ_t ?

Mean field for the bosons: The Hartree equation

$$H = \sum_{j=1}^N -\Delta_j + \sum_{j=1}^N A_t(x_j) + (N-1)^{-1} \sum_{k < j} V(x_j - x_k)$$

Interaction „felt“ by each particle of order one

$$\Psi_0 = \prod_{j=1}^N \phi_0(x_j) \quad id_t \Psi_t = H_t \Psi_t$$

Interaction destroys product structure.

Question:

- ▶ In which regimes: $\Psi_t \approx \prod_{j=1}^N \phi_t(x_j)$ (in what sense?)
- ▶ What is ϕ_t ?

Mean field for the bosons: The Hartree equation

$$H = \sum_{j=1}^N -\Delta_j + \sum_{j=1}^N A_t(x_j) + (N-1)^{-1} \sum_{k < j} V(x_j - x_k)$$

Interaction „felt“ by each particle of order one

$$\Psi_0 = \prod_{j=1}^N \phi_0(x_j) \quad id_t \Psi_t = H_t \Psi_t$$

Interaction destroys product structure.

Question:

- ▶ In which regimes: $\Psi_t \approx \prod_{j=1}^N \phi_t(x_j)$ (in what sense?)
- ▶ What is ϕ_t ?

Mean field for the bosons: The Hartree equation

$$H = \sum_{j=1}^N -\Delta_j + \sum_{j=1}^N A_t(x_j) + (N-1)^{-1} \sum_{k < j} V(x_j - x_k)$$

Interaction „felt“ by each particle of order one

$$\Psi_0 = \prod_{j=1}^N \phi_0(x_j) \quad id_t \Psi_t = H_t \Psi_t$$

Interaction destroys product structure.

Question:

- ▶ In which regimes: $\Psi_t \approx \prod_{j=1}^N \phi_t(x_j)$ (in what sense?)
- ▶ What is ϕ_t ?

Mean field for the bosons: The Hartree equation

$$H = \sum_{j=1}^N -\Delta_j + \sum_{j=1}^N A_t(x_j) + (N-1)^{-1} \sum_{k < j} V(x_j - x_k)$$

Interaction „felt“ by each particle of order one

$$\Psi_0 = \prod_{j=1}^N \phi_0(x_j) \quad id_t \Psi_t = H_t \Psi_t$$

Interaction destroys product structure.

Question:

- ▶ In which regimes: $\Psi_t \approx \prod_{j=1}^N \phi_t(x_j)$ (in what sense?)
- ▶ What is ϕ_t ?

Mean field for the bosons: The Hartree equation

- ▶ Assuming $\Psi_t \approx \prod_{j=1}^N \phi_t(x_j)$ and finding ϕ_t : easy.
- ▶ Proving $\Psi_t \approx \prod_{j=1}^N \phi_t(x_j)$: hard.
In particular error propagation
- ▶ Example: Gross-Pitaevskii for dilute gases: NOT a mean field situation!

Mean field for the bosons: The Hartree equation

- ▶ Assuming $\Psi_t \approx \prod_{j=1}^N \phi_t(x_j)$ and finding ϕ_t : easy.
- ▶ Proving $\Psi_t \approx \prod_{j=1}^N \phi_t(x_j)$: hard.
In particular error propagation
- ▶ Example: Gross-Pitaevskii for dilute gases: NOT a mean field situation!

Mean field for the bosons: The Hartree equation

- ▶ Assuming $\Psi_t \approx \prod_{j=1}^N \phi_t(x_j)$ and finding ϕ_t : easy.
- ▶ Proving $\Psi_t \approx \prod_{j=1}^N \phi_t(x_j)$: hard.
In particular error propagation
- ▶ Example: Gross-Pitaevskii for dilute gases: NOT a mean field situation!

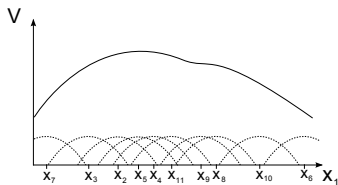
Mean field for the bosons: The Hartree equation

- ▶ Assuming $\Psi_t \approx \prod_{j=1}^N \phi_t(x_j)$ and finding ϕ_t : easy.
- ▶ Proving $\Psi_t \approx \prod_{j=1}^N \phi_t(x_j)$: hard.
In particular error propagation
- ▶ Example: Gross-Pitaevskii for dilute gases: NOT a mean field situation!

Mean field for the bosons: The Hartree equation

- ▶ Assuming $\Psi_t \approx \prod_{j=1}^N \phi_t(x_j)$ and finding ϕ_t : easy.
- ▶ Proving $\Psi_t \approx \prod_{j=1}^N \phi_t(x_j)$: hard.
In particular error propagation
- ▶ Example: Gross-Pitaevskii for dilute gases: NOT a mean field situation!

Mean field for “particle 1”



$W(x_1) = (N - 1)^{-1} \sum_{j=2}^N V(x_1 - x_j)$ for fixed, $|\phi_0|^2$ - distributed x_2, \dots, x_N .

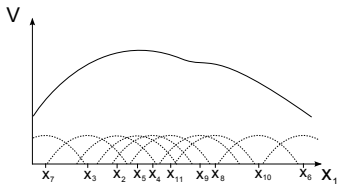
Law of large numbers: $|\phi_0|^2$ close to the empirical density ρ_0 .

$W(x_1) \approx V \star |\phi_0|^2(x_1)$ (“Mean field”).

Effective Dynamics: Hartree equation

$$id_t \phi_t = (-\Delta + A_t + V \star |\phi_t|^2) \phi_t .$$

Mean field for “particle 1”



$W(x_1) = (N - 1)^{-1} \sum_{j=2}^N V(x_1 - x_j)$ for fixed, $|\phi_0|^2$ - distributed x_2, \dots, x_N .

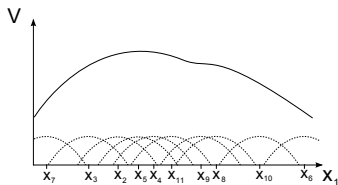
Law of large numbers: $|\phi_0|^2$ close to the empirical density ρ_0 .

$W(x_1) \approx V \star |\phi_0|^2(x_1)$ (“Mean field”).

Effective Dynamics: Hartree equation

$$id_t \phi_t = (-\Delta + A_t + V \star |\phi_t|^2) \phi_t .$$

Mean field for “particle 1”



$W(x_1) = (N - 1)^{-1} \sum_{j=2}^N V(x_1 - x_j)$ for fixed, $|\phi_0|^2$ - distributed x_2, \dots, x_N .

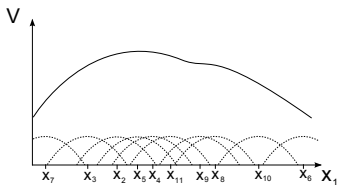
Law of large numbers: $|\phi_0|^2$ close to the empirical density ρ_0 .

$W(x_1) \approx V \star |\phi_0|^2(x_1)$ (“Mean field”).

Effective Dynamics: Hartree equation

$$id_t \phi_t = (-\Delta + A_t + V \star |\phi_t|^2) \phi_t .$$

Mean field for “particle 1”



$W(x_1) = (N - 1)^{-1} \sum_{j=2}^N V(x_1 - x_j)$ for fixed, $|\phi_0|^2$ - distributed x_2, \dots, x_N .

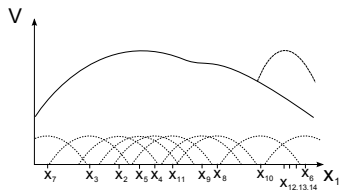
Law of large numbers: $|\phi_0|^2$ close to the empirical density ρ_0 .

$W(x_1) \approx V \star |\phi_0|^2(x_1)$ (“Mean field”).

Effective Dynamics: Hartree equation

$$id_t \phi_t = (-\Delta + A_t + V \star |\phi_t|^2) \phi_t .$$

Grönwall argument

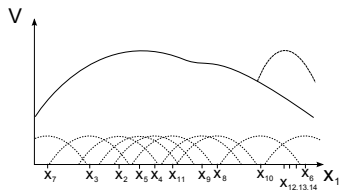


Let α_t be a measure for the dirt in the condensate:

$$d_t \alpha_t \leq C(\alpha_t + o(1))$$

Grönwall: α_t stays small if α_0 was small ($\alpha_t \leq e^{Ct} \alpha_0 + o(1)$)

Grönwall argument

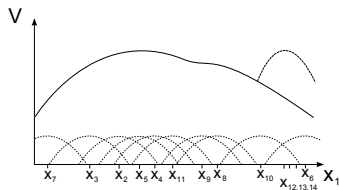


Let α_t be a measure for the dirt in the condensate:

$$d_t \alpha_t \leq C(\alpha_t + o(1))$$

Grönwall: α_t stays small if α_0 was small ($\alpha_t \leq e^{Ct} \alpha_0 + o(1)$)

Grönwall argument



Let α_t be a measure for the dirt in the condensate:

$$d_t \alpha_t \leq C(\alpha_t + \mathcal{O}(1))$$

Grönwall: α_t stays small if α_0 was small ($\alpha_t \leq e^{Ct} \alpha_0 + \mathcal{O}(1)$)

Usual approach

Control reduced one particle density matrix $\mu_1^{\Psi_t}$

$d_t \mu_1^{\Psi_t}$ depends on $\mu_2^{\Psi_t}$

$d_t \mu_2^{\Psi_t}$ depends on $\mu_3^{\Psi_t}$

...

Technically difficult (e.g. Erdős, Schlein, Yau (2009))

Usual approach

Control reduced one particle density matrix $\mu_1^{\Psi_t}$

$d_t \mu_1^{\Psi_t}$ depends on $\mu_2^{\Psi_t}$

$d_t \mu_2^{\Psi_t}$ depends on $\mu_3^{\Psi_t}$

...

Technically difficult (e.g. Erdős, Schlein, Yau (2009))

Usual approach

Control reduced one particle density matrix $\mu_1^{\Psi_t}$

$d_t \mu_1^{\Psi_t}$ depends on $\mu_2^{\Psi_t}$

$d_t \mu_2^{\Psi_t}$ depends on $\mu_3^{\Psi_t}$

...

Technically difficult (e.g. Erdős, Schlein, Yau (2009))

Usual approach

Control reduced one particle density matrix $\mu_1^{\Psi_t}$

$d_t \mu_1^{\Psi_t}$ depends on $\mu_2^{\Psi_t}$

$d_t \mu_2^{\Psi_t}$ depends on $\mu_3^{\Psi_t}$

...

Technically difficult (e.g. Erdős, Schlein, Yau (2009))

Usual approach

Control reduced one particle density matrix $\mu_1^{\Psi_t}$

$d_t \mu_1^{\Psi_t}$ depends on $\mu_2^{\Psi_t}$

$d_t \mu_2^{\Psi_t}$ depends on $\mu_3^{\Psi_t}$

...

Technically difficult (e.g. Erdős, Schlein, Yau (2009))

Usual approach

Control reduced one particle density matrix $\mu_1^{\Psi_t}$

$d_t \mu_1^{\Psi_t}$ depends on $\mu_2^{\Psi_t}$

$d_t \mu_2^{\Psi_t}$ depends on $\mu_3^{\Psi_t}$

...

Technically difficult (e.g. Erdős, Schlein, Yau (2009))

Introduction of a counting measure:

For every $j = 1, \dots, N$ and $\phi \in L^2(\mathbb{R}^3)$ let p_j^ϕ be the projector given by

$$p_j^\phi = |\phi\rangle\langle\phi|_j.$$

Let $q_j^\phi = 1 - p_j^\phi$.

According to Ψ "expected relative number" of particles not in the state ϕ

$$\alpha(\Psi, \phi) = N^{-1} \sum_{j=1}^N \langle \Psi, q_j^\phi \Psi \rangle = \langle \Psi, q_k^\phi \Psi \rangle = \|q_k^\phi \Psi\|^2.$$

$$\alpha(\Psi, \phi) = 0 \quad \Leftrightarrow \quad \Psi = \prod_{j=1}^N \phi(x_j).$$

Introduction of a counting measure:

For every $j = 1, \dots, N$ and $\phi \in L^2(\mathbb{R}^3)$ let p_j^ϕ be the projector given by

$$p_j^\phi = |\phi\rangle\langle\phi|_j.$$

Let $q_j^\phi = 1 - p_j^\phi$.

According to Ψ "expected relative number" of particles not in the state ϕ

$$\alpha(\Psi, \phi) = N^{-1} \sum_{j=1}^N \langle \Psi, q_j^\phi \Psi \rangle = \langle \Psi, q_k^\phi \Psi \rangle = \|q_k^\phi \Psi\|^2.$$

$$\alpha(\Psi, \phi) = 0 \quad \Leftrightarrow \quad \Psi = \prod_{j=1}^N \phi(x_j).$$

Introduction of a counting measure:

For every $j = 1, \dots, N$ and $\phi \in L^2(\mathbb{R}^3)$ let p_j^ϕ be the projector given by

$$p_j^\phi = |\phi\rangle\langle\phi|_j.$$

Let $q_j^\phi = 1 - p_j^\phi$.

According to Ψ "expected relative number" of particles not in the state ϕ

$$\alpha(\Psi, \phi) = N^{-1} \sum_{j=1}^N \langle \Psi, q_j^\phi \Psi \rangle = \langle \Psi, q_k^\phi \Psi \rangle = \|q_k^\phi \Psi\|^2.$$

$$\alpha(\Psi, \phi) = 0 \quad \Leftrightarrow \quad \Psi = \prod_{j=1}^N \phi(x_j).$$

Introduction of a counting measure:

For every $j = 1, \dots, N$ and $\phi \in L^2(\mathbb{R}^3)$ let p_j^ϕ be the projector given by

$$p_j^\phi = |\phi\rangle\langle\phi|_j.$$

Let $q_j^\phi = 1 - p_j^\phi$.

According to Ψ “expected relative number” of particles not in the state ϕ

$$\alpha(\Psi, \phi) = N^{-1} \sum_{j=1}^N \langle \Psi, q_j^\phi \Psi \rangle = \langle \Psi, q_k^\phi \Psi \rangle = \|q_k^\phi \Psi\|^2.$$

$$\alpha(\Psi, \phi) = 0 \quad \Leftrightarrow \quad \Psi = \prod_{j=1}^N \phi(x_j).$$

Introduction of a counting measure:

For every $j = 1, \dots, N$ and $\phi \in L^2(\mathbb{R}^3)$ let p_j^ϕ be the projector given by

$$p_j^\phi = |\phi\rangle\langle\phi|_j.$$

Let $q_j^\phi = 1 - p_j^\phi$.

According to Ψ “expected relative number” of particles not in the state ϕ

$$\alpha(\Psi, \phi) = N^{-1} \sum_{j=1}^N \langle \Psi, q_j^\phi \Psi \rangle = \langle \Psi, q_k^\phi \Psi \rangle = \|q_k^\phi \Psi\|^2.$$

$$\alpha(\Psi, \phi) = 0 \quad \Leftrightarrow \quad \Psi = \prod_{j=1}^N \phi(x_j).$$

Example:

$$\Psi(x_1, \dots, x_N) = \left(\chi(x_1, \dots, x_k) \prod_{j=k+1}^N \phi(x_j) \right)_{sym}$$

with $\chi \in L^2(\mathbb{R}^{3k})$, $p_j \chi = 0$ for all $1 \leq j \leq k$

$$\Rightarrow \alpha(\Psi, \phi) = k/N.$$

Lemma: $\alpha(\Psi_t, \phi_t) \rightarrow 0$ is equivalent to convergence of μ_t to $|\phi_t\rangle\langle\phi_t|$ in operator norm.

Example:

$$\Psi(x_1, \dots, x_N) = \left(\chi(x_1, \dots, x_k) \prod_{j=k+1}^N \phi(x_j) \right)_{sym}$$

with $\chi \in L^2(\mathbb{R}^{3k})$, $p_j \chi = 0$ for all $1 \leq j \leq k$

$$\Rightarrow \alpha(\Psi, \phi) = k/N.$$

Lemma: $\alpha(\Psi_t, \phi_t) \rightarrow 0$ is equivalent to convergence of μ_t to $|\phi_t\rangle\langle\phi_t|$ in operator norm.

Example:

$$\Psi(x_1, \dots, x_N) = \left(\chi(x_1, \dots, x_k) \prod_{j=k+1}^N \phi(x_j) \right)_{sym}$$

with $\chi \in L^2(\mathbb{R}^{3k})$, $p_j \chi = 0$ for all $1 \leq j \leq k$

$$\Rightarrow \alpha(\Psi, \phi) = k/N.$$

Lemma: $\alpha(\Psi_t, \phi_t) \rightarrow 0$ is equivalent to convergence of μ_t to $|\phi_t\rangle\langle\phi_t|$ in operator norm.

Deriving the mean field equation:

Goal: show that $|d_t \alpha_t| < C(\alpha_t + o(1)) \Rightarrow \alpha_t \ll 1$ by Grönwall

Let

$$h_j := -\Delta_j + A + V \star |\phi_t|^2(x_j) .$$

Since

$$d_t q_j^{\phi_t} = -i[h_j, q_j^{\phi_t}]$$

$$d_t \alpha(\Psi_t, \phi_t) = i \langle \Psi_t, [H - h_1, q_1^{\phi_t}] \Psi_t \rangle .$$

Deriving the mean field equation:

Goal: show that $|d_t \alpha_t| < C(\alpha_t + o(1)) \Rightarrow \alpha_t \ll 1$ by Grönwall

Let

$$h_j := -\Delta_j + A + V \star |\phi_t|^2(x_j) .$$

Since

$$d_t q_j^{\phi_t} = -i[h_j, q_j^{\phi_t}]$$

$$d_t \alpha(\Psi_t, \phi_t) = i \langle \Psi_t, [H - h_1, q_1^{\phi_t}] \Psi_t \rangle .$$

Deriving the mean field equation:

Goal: show that $|d_t \alpha_t| < C(\alpha_t + o(1)) \Rightarrow \alpha_t \ll 1$ by Grönwall

Let

$$h_j := -\Delta_j + A + V \star |\phi_t|^2(x_j) .$$

Since

$$d_t q_j^{\phi_t} = -i[h_j, q_j^{\phi_t}]$$

$$d_t \alpha(\Psi_t, \phi_t) = i \langle \Psi_t, [H - h_1, q_1^{\phi_t}] \Psi_t \rangle .$$

Deriving the mean field equation:

Goal: show that $|d_t \alpha_t| < C(\alpha_t + o(1)) \Rightarrow \alpha_t \ll 1$ by Grönwall

Let

$$h_j := -\Delta_j + A + V \star |\phi_t|^2(x_j) .$$

Since

$$d_t q_j^{\phi_t} = -i[h_j, q_j^{\phi_t}]$$

$$d_t \alpha(\Psi_t, \phi_t) = i \langle \Psi_t, [H - h_1, q_1^{\phi_t}] \Psi_t \rangle .$$

Deriving the mean field equation:

Goal: show that $|d_t \alpha_t| < C(\alpha_t + o(1)) \Rightarrow \alpha_t \ll 1$ by Grönwall

Let

$$h_j := -\Delta_j + A + V \star |\phi_t|^2(x_j) .$$

Since

$$d_t q_j^{\phi_t} = -i[h_j, q_j^{\phi_t}]$$

$$d_t \alpha(\Psi_t, \phi_t) = i \langle \Psi_t, [H - h_1, q_1^{\phi_t}] \Psi_t \rangle .$$

$$d_t \alpha(\Psi_t, \phi_t) = i \langle \Psi_t, [-\Delta_1 + (N-1)^{-1} \sum_{j=2}^N V(x_1 - x_j) - h_1, q_1^{\phi_t}] \Psi_t \rangle$$

$$= i \langle \Psi_t, [V(x_1 - x_2) - V \star |\phi_t|^2(x_1), q_1^{\phi_t}] \Psi_t \rangle$$

$$= i (\langle \Psi_t, (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} \Psi_t \rangle - c.c.)$$

$$= i (\langle \Psi_t, (p_1^{\phi_t} + q_1^{\phi_t})(p_2^{\phi_t} + q_2^{\phi_t})(V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t}(p_2^{\phi_t} + q_2^{\phi_t}) \Psi_t \rangle - c.c.) .$$

All terms $\langle \Psi_t, A_{12} \Psi_t \rangle$ where A_{12} is either selfadjoint or invariant under adjunction plus simultaneous exchange of the variables x_1 and x_2 cancel out.

$$d_t \alpha(\Psi_t, \phi_t) = i \langle \Psi_t, [-\Delta_1 + (N-1)^{-1} \sum_{j=2}^N V(x_1 - x_j) - h_1, q_1^{\phi_t}] \Psi_t \rangle$$

$$= i \langle \Psi_t, [V(x_1 - x_2) - V \star |\phi_t|^2(x_1), q_1^{\phi_t}] \Psi_t \rangle$$

$$= i (\langle \Psi_t, (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} \Psi_t \rangle - c.c.)$$

$$= i (\langle \Psi_t, (p_1^{\phi_t} + q_1^{\phi_t})(p_2^{\phi_t} + q_2^{\phi_t}) (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} (p_2^{\phi_t} + q_2^{\phi_t}) \Psi_t \rangle - c.c.) .$$

All terms $\langle \Psi_t, A_{12} \Psi_t \rangle$ where A_{12} is either selfadjoint or invariant under adjunction plus simultaneous exchange of the variables x_1 and x_2 cancel out.

$$d_t \alpha(\Psi_t, \phi_t) = i \langle \Psi_t, [-\Delta_1 + (N-1)^{-1} \sum_{j=2}^N V(x_1 - x_j) - h_1, q_1^{\phi_t}] \Psi_t \rangle$$

$$= i \langle \Psi_t, [V(x_1 - x_2) - V \star |\phi_t|^2(x_1), q_1^{\phi_t}] \Psi_t \rangle$$

$$= i (\langle \Psi_t, (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} \Psi_t \rangle - c.c.)$$

$$= i (\langle \Psi_t, (p_1^{\phi_t} + q_1^{\phi_t})(p_2^{\phi_t} + q_2^{\phi_t})(V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t}(p_2^{\phi_t} + q_2^{\phi_t}) \Psi_t \rangle - c.c.) .$$

All terms $\langle \Psi_t, A_{12} \Psi_t \rangle$ where A_{12} is either selfadjoint or invariant under adjunction plus simultaneous exchange of the variables x_1 and x_2 cancel out.

$$d_t \alpha(\Psi_t, \phi_t) = i \langle \Psi_t, [-\Delta_1 + (N-1)^{-1} \sum_{j=2}^N V(x_1 - x_j) - h_1, q_1^{\phi_t}] \Psi_t \rangle$$

$$= i \langle \Psi_t, [V(x_1 - x_2) - V \star |\phi_t|^2(x_1), q_1^{\phi_t}] \Psi_t \rangle$$

$$= i (\langle \Psi_t, (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} \Psi_t \rangle - c.c.)$$

$$= i (\langle \Psi_t, (p_1^{\phi_t} + q_1^{\phi_t})(p_2^{\phi_t} + q_2^{\phi_t}) (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} (p_2^{\phi_t} + q_2^{\phi_t}) \Psi_t \rangle - c.c.) .$$

All terms $\langle \Psi_t, A_{12} \Psi_t \rangle$ where A_{12} is either selfadjoint or invariant under adjunction plus simultaneous exchange of the variables x_1 and x_2 cancel out.

$$d_t \alpha(\Psi_t, \phi_t) = i \langle \Psi_t, [-\Delta_1 + (N-1)^{-1} \sum_{j=2}^N V(x_1 - x_j) - h_1, q_1^{\phi_t}] \Psi_t \rangle$$

$$= i \langle \Psi_t, [V(x_1 - x_2) - V \star |\phi_t|^2(x_1), q_1^{\phi_t}] \Psi_t \rangle$$

$$= i (\langle \Psi_t, (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} \Psi_t \rangle - c.c.)$$

$$= i (\langle \Psi_t, (p_1^{\phi_t} + q_1^{\phi_t})(p_2^{\phi_t} + q_2^{\phi_t})(V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t}(p_2^{\phi_t} + q_2^{\phi_t}) \Psi_t \rangle - c.c.) .$$

All terms $\langle \Psi_t, A_{12} \Psi_t \rangle$ where A_{12} is either selfadjoint or invariant under adjunction plus simultaneous exchange of the variables x_1 and x_2 cancel out.

Only three types remain: $pp \dots pq$, $pp \dots qq$ and $pq \dots qq$

$$\begin{aligned}
 I &= i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} p_2^{\phi_t} \Psi_t \rangle \\
 &= i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle - i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V \star |\phi_t|^2(x_1) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle
 \end{aligned}$$

$$II = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle$$

$$III = i \langle \Psi_t, p_1^{\phi_t} q_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle .$$

Only three types remain: $pp \dots pq$, $pp \dots qq$ and $pq \dots qq$

$$I = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} p_2^{\phi_t} \Psi_t \rangle$$

$$= i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle - i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V \star |\phi_t|^2(x_1) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle$$

$$II = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle$$

$$III = i \langle \Psi_t, p_1^{\phi_t} q_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle .$$

Only three types remain: $pp \dots pq$, $pp \dots qq$ and $pq \dots qq$

$$I = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} p_2^{\phi_t} \Psi_t \rangle$$

$$= i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle - i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V \star |\phi_t|^2(x_1) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle$$

$$II = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle$$

$$III = i \langle \Psi_t, p_1^{\phi_t} q_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle .$$

Only three types remain: $pp \dots pq$, $pp \dots qq$ and $pq \dots qq$

$$I = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} p_2^{\phi_t} \Psi_t \rangle$$

$$= i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle - i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V \star |\phi_t|^2(x_1) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle$$

$$II = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle$$

$$III = i \langle \Psi_t, p_1^{\phi_t} q_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle .$$

Only three types remain: $pp \dots pq$, $pp \dots qq$ and $pq \dots qq$

$$\begin{aligned}
 I &= i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} p_2^{\phi_t} \Psi_t \rangle \\
 &= i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle - i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V \star |\phi_t|^2(x_1) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle
 \end{aligned}$$

$$II = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle$$

$$III = i \langle \Psi_t, p_1^{\phi_t} q_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle .$$

Only three types remain: $pp \dots pq$, $pp \dots qq$ and $pq \dots qq$

$$\begin{aligned}
 I &= i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} p_2^{\phi_t} \Psi_t \rangle \\
 &= i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle - i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V \star |\phi_t|^2(x_1) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle
 \end{aligned}$$

$$II = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle$$

$$III = i \langle \Psi_t, p_1^{\phi_t} q_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle .$$

$$I : p_2^{\phi_t} V(x_1 - x_2) p_2^{\phi_t} =$$

Only three types remain: $pp \dots pq$, $pp \dots qq$ and $pq \dots qq$

$$I = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} p_2^{\phi_t} \Psi_t \rangle$$

$$= i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle - i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V \star |\phi_t|^2(x_1) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle$$

$$II = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle$$

$$III = i \langle \Psi_t, p_1^{\phi_t} q_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle .$$

$$I : p_2^{\phi_t} V(x_1 - x_2) p_2^{\phi_t} = |\phi_t(x_2)\rangle \langle \phi_t(x_2)| V(x_1 - x_2) |\phi_t(x_2)\rangle \langle \phi_t(x_2)|$$

Only three types remain: $pp \dots pq$, $pp \dots qq$ and $pq \dots qq$

$$I = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} p_2^{\phi_t} \Psi_t \rangle$$

$$= i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle - i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V \star |\phi_t|^2(x_1) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle$$

$$II = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle$$

$$III = i \langle \Psi_t, p_1^{\phi_t} q_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle .$$

$$I : p_2^{\phi_t} V(x_1 - x_2) p_2^{\phi_t} = |\phi_t(x_2)\rangle \langle \phi_t(x_2)| V(x_1 - x_2) |\phi_t(x_2)\rangle \langle \phi_t(x_2)|$$

$$= |\phi_t(x_2)\rangle V \star |\phi_t|^2(x_1) \langle \phi_t(x_2)| = p_2^{\phi_t} V \star |\phi_t|^2$$

$$= p_2^{\phi_t} V \star |\phi_t|^2 p_2^{\phi_t}$$

$$\Rightarrow I = 0.$$

Only three types remain: $pp \dots pq$, $pp \dots qq$ and $pq \dots qq$

$$I = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} p_2^{\phi_t} \Psi_t \rangle$$

$$= i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle - i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V \star |\phi_t|^2(x_1) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle$$

$$II = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle$$

$$III = i \langle \Psi_t, p_1^{\phi_t} q_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle .$$

$$I : p_2^{\phi_t} V(x_1 - x_2) p_2^{\phi_t} = |\phi_t(x_2)\rangle \langle \phi_t(x_2)| V(x_1 - x_2) |\phi_t(x_2)\rangle \langle \phi_t(x_2)|$$

$$= |\phi_t(x_2)\rangle V \star |\phi_t|^2(x_1) \langle \phi_t(x_2)| = p_2^{\phi_t} V \star |\phi_t|^2$$

$$= p_2^{\phi_t} V \star |\phi_t|^2 p_2^{\phi_t}$$

$$\Rightarrow I = 0.$$

Only three types remain: $pp \dots pq$, $pp \dots qq$ and $pq \dots qq$

$$I = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} p_2^{\phi_t} \Psi_t \rangle$$

$$= i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle - i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V \star |\phi_t|^2(x_1) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle$$

$$II = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle$$

$$III = i \langle \Psi_t, p_1^{\phi_t} q_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle .$$

$$I : p_2^{\phi_t} V(x_1 - x_2) p_2^{\phi_t} = |\phi_t(x_2)\rangle \langle \phi_t(x_2)| V(x_1 - x_2) |\phi_t(x_2)\rangle \langle \phi_t(x_2)|$$

$$= |\phi_t(x_2)\rangle V \star |\phi_t|^2(x_1) \langle \phi_t(x_2)| = p_2^{\phi_t} V \star |\phi_t|^2$$

$$= p_2^{\phi_t} V \star |\phi_t|^2 p_2^{\phi_t}$$

$\Rightarrow I = 0.$

Only three types remain: $pp \dots pq$, $pp \dots qq$ and $pq \dots qq$

$$I = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} p_2^{\phi_t} \Psi_t \rangle$$

$$= i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle - i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V \star |\phi_t|^2(x_1) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle$$

$$II = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle$$

$$III = i \langle \Psi_t, p_1^{\phi_t} q_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle .$$

$$I : p_2^{\phi_t} V(x_1 - x_2) p_2^{\phi_t} = |\phi_t(x_2)\rangle \langle \phi_t(x_2)| V(x_1 - x_2) |\phi_t(x_2)\rangle \langle \phi_t(x_2)|$$

$$= |\phi_t(x_2)\rangle V \star |\phi_t|^2(x_1) \langle \phi_t(x_2)| = p_2^{\phi_t} V \star |\phi_t|^2$$

$$= p_2^{\phi_t} V \star |\phi_t|^2 p_2^{\phi_t}$$

$\Rightarrow I = 0.$

Only three types remain: $pp \dots pq$, $pp \dots qq$ and $pq \dots qq$

$$I = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} p_2^{\phi_t} \Psi_t \rangle$$

$$= i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle - i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V \star |\phi_t|^2(x_1) p_2^{\phi_t} q_1^{\phi_t} \Psi_t \rangle$$

$$II = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle$$

$$III = i \langle \Psi_t, p_1^{\phi_t} q_2^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle .$$

$$I : p_2^{\phi_t} V(x_1 - x_2) p_2^{\phi_t} = |\phi_t(x_2)\rangle \langle \phi_t(x_2)| V(x_1 - x_2) |\phi_t(x_2)\rangle \langle \phi_t(x_2)|$$

$$= |\phi_t(x_2)\rangle V \star |\phi_t|^2(x_1) \langle \phi_t(x_2)| = p_2^{\phi_t} V \star |\phi_t|^2$$

$$= p_2^{\phi_t} V \star |\phi_t|^2 p_2^{\phi_t}$$

$$\Rightarrow I = 0.$$

Only three types remain: $pp \dots pq$, $pp \dots qq$ and $pq \dots qq$

$$I = 0$$

$$II = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle$$

$$III = i \langle \Psi_t, q_2^{\phi_t} p_1^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle .$$

$$|III| \leq \|q_2^{\phi_t} \Psi_t\|^2 \|V\|_{\infty} = C\alpha(\Psi_t, \phi_t)$$

$$d_t \alpha(\Psi_t, \phi_t) \leq C(\alpha(\Psi_t, \phi_t) + N^{-1})$$

Only three types remain: $pp \dots pq$, $pp \dots qq$ and $pq \dots qq$

$$I = 0$$

$$II = i \langle \Psi_t, p_1^{\phi_t} p_2^{\phi_t} V(x_1 - x_2) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle$$

$$III = i \langle \Psi_t, q_2^{\phi_t} p_1^{\phi_t} (V(x_1 - x_2) - V \star |\phi_t|^2(x_1)) q_1^{\phi_t} q_2^{\phi_t} \Psi_t \rangle .$$

$$|III| \leq \|q_2^{\phi_t} \Psi_t\|^2 \|V\|_{\infty} = C\alpha(\Psi_t, \phi_t)$$

$$d_t \alpha(\Psi_t, \phi_t) \leq C(\alpha(\Psi_t, \phi_t) + N^{-1})$$

Fermions

- ▶ Microscopic system: $\Psi_0 = \Lambda_{j=1}^N \phi_0^j$

$$H = \sum_{j=1}^N -\Delta_j + \sum_{j=1}^N A_t(x_j) + N^{-2/3} \sum_{k < j} V(x_j - x_k)$$

$N^{-2/3}$ -coupling: volume N , Coulomb interaction

$$\int_0^{N^{1/3}} x^{-1} d^3x \sim N^{2/3}$$

- ▶ Macroscopic equation: Hartree (Fock)

$$d_t \phi_t^j = \left(-\Delta + A + N^{-1} \sum_{k=1}^N |\phi_t^k|^2 \star V \right) \phi_t^j$$

$$- \left(N^{-1} \sum_{k=1}^N \left(\phi_t^{k,*} \phi_t^j \right) \star V \right) \phi_t^k$$

- ▶ $\rho = \sum_{j=1}^N |\phi_t^j\rangle \langle \phi_t^j|$

- ▶ Difficulties: $I \neq 0$

$N^{-2/3}$ instead of N^{-1}

Fermions

- ▶ Microscopic system: $\Psi_0 = \Lambda_{j=1}^N \phi_0^j$

$$H = \sum_{j=1}^N -\Delta_j + \sum_{j=1}^N A_t(x_j) + N^{-2/3} \sum_{k < j} V(x_j - x_k)$$

$N^{-2/3}$ -coupling: volume N , Coulomb interaction

$$\int_0^{N^{1/3}} x^{-1} d^3x \sim N^{2/3}$$

- ▶ Macroscopic equation: Hartree (Fock)

$$d_t \phi_t^j = \left(-\Delta + A + N^{-1} \sum_{k=1}^N |\phi_t^k|^2 \star V \right) \phi_t^j$$

$$- \left(N^{-1} \sum_{k=1}^N \left(\phi_t^{k*} \phi_t^j \right) \star V \right) \phi_t^k$$

- ▶ $\rho = \sum_{j=1}^N |\phi_t^j\rangle \langle \phi_t^j|$

- ▶ Difficulties: $I \neq 0$

$N^{-2/3}$ instead of N^{-1}

Fermions

- ▶ Microscopic system: $\Psi_0 = \Lambda_{j=1}^N \phi_0^j$
 $H = \sum_{j=1}^N -\Delta_j + \sum_{j=1}^N A_t(x_j) + N^{-2/3} \sum_{k < j} V(x_j - x_k)$
 $N^{-2/3}$ -coupling: volume N , Coulomb interaction

$$\int_0^{N^{1/3}} x^{-1} d^3x \sim N^{2/3}$$

- ▶ Macroscopic equation: Hartree (Fock)
 $d_t \phi_t^j = \left(-\Delta + A + N^{-1} \sum_{k=1}^N |\phi_t^k|^2 * V \right) \phi_t^j$
 $- \left(N^{-1} \sum_{k=1}^N \left(\phi_t^{k*} \phi_t^j \right) * V \right) \phi_t^k$
- ▶ $\rho = \sum_{j=1}^N |\phi_t^j\rangle \langle \phi_t^j|$
- ▶ Difficulties: $I \neq 0$
 $N^{-2/3}$ instead of N^{-1}

Fermions

- ▶ Microscopic system: $\Psi_0 = \Lambda_{j=1}^N \phi_0^j$
 $H = \sum_{j=1}^N -\Delta_j + \sum_{j=1}^N A_t(x_j) + N^{-2/3} \sum_{k < j} V(x_j - x_k)$
 $N^{-2/3}$ -coupling: volume N , Coulomb interaction

$$\int_0^{N^{1/3}} x^{-1} d^3x \sim N^{2/3}$$

- ▶ Macroscopic equation: Hartree (Fock)
 $d_t \phi_t^j = \left(-\Delta + A + N^{-1} \sum_{k=1}^N |\phi_t^k|^2 * V \right) \phi_t^j$
 $- \left(N^{-1} \sum_{k=1}^N \left(\phi_t^{k*} \phi_t^j \right) * V \right) \phi_t^k$
- ▶ $\rho = \sum_{j=1}^N |\phi_t^j\rangle \langle \phi_t^j|$
- ▶ Difficulties: $I \neq 0$
 $N^{-2/3}$ instead of N^{-1}

Fermions

- ▶ Microscopic system: $\Psi_0 = \Lambda_{j=1}^N \phi_0^j$
 $H = \sum_{j=1}^N -\Delta_j + \sum_{j=1}^N A_t(x_j) + N^{-2/3} \sum_{k < j} V(x_j - x_k)$
 $N^{-2/3}$ -coupling: volume N , Coulomb interaction

$$\int_0^{N^{1/3}} x^{-1} d^3x \sim N^{2/3}$$

- ▶ Macroscopic equation: Hartree (Fock)
 $d_t \phi_t^j = \left(-\Delta + A + N^{-1} \sum_{k=1}^N |\phi_t^k|^2 * V \right) \phi_t^j$
 $- \left(N^{-1} \sum_{k=1}^N \left(\phi_t^{k*} \phi_t^j \right) * V \right) \phi_t^k$
- ▶ $\rho = \sum_{j=1}^N |\phi_t^j\rangle \langle \phi_t^j|$
- ▶ Difficulties: $I \neq 0$
 $N^{-2/3}$ instead of N^{-1}

Fermions

- ▶ Microscopic system: $\Psi_0 = \Lambda_{j=1}^N \phi_0^j$
 $H = \sum_{j=1}^N -\Delta_j + \sum_{j=1}^N A_t(x_j) + N^{-2/3} \sum_{k < j} V(x_j - x_k)$
 $N^{-2/3}$ -coupling: volume N , Coulomb interaction

$$\int_0^{N^{1/3}} x^{-1} d^3x \sim N^{2/3}$$

- ▶ Macroscopic equation: Hartree (Fock)
 $d_t \phi_t^j = \left(-\Delta + A + N^{-1} \sum_{k=1}^N |\phi_t^k|^2 \star V \right) \phi_t^j$
 $- \left(N^{-1} \sum_{k=1}^N \left(\phi_t^{k,*} \phi_t^j \right) \star V \right) \phi_t^k$
- ▶ $\rho = \sum_{j=1}^N |\phi_t^j\rangle \langle \phi_t^j|$
- ▶ Difficulties: $I \neq 0$
 $N^{-2/3}$ instead of N^{-1}

Fermions

- ▶ Microscopic system: $\Psi_0 = \Lambda_{j=1}^N \phi_0^j$
 $H = \sum_{j=1}^N -\Delta_j + \sum_{j=1}^N A_t(x_j) + N^{-2/3} \sum_{k < j} V(x_j - x_k)$
 $N^{-2/3}$ -coupling: volume N , Coulomb interaction

$$\int_0^{N^{1/3}} x^{-1} d^3x \sim N^{2/3}$$

- ▶ Macroscopic equation: Hartree (Fock)
 $d_t \phi_t^j = \left(-\Delta + A + N^{-1} \sum_{k=1}^N |\phi_t^k|^2 \star V \right) \phi_t^j$
 $- \left(N^{-1} \sum_{k=1}^N \left(\phi_t^{k,*} \phi_t^j \right) \star V \right) \phi_t^k$

- ▶ $\rho = \sum_{j=1}^N |\phi_t^j\rangle \langle \phi_t^j|$

- ▶ Difficulties: $l \neq 0$
 $N^{-2/3}$ instead of N^{-1}

Fermions

- ▶ Microscopic system: $\Psi_0 = \Lambda_{j=1}^N \phi_0^j$
 $H = \sum_{j=1}^N -\Delta_j + \sum_{j=1}^N A_t(x_j) + N^{-2/3} \sum_{k < j} V(x_j - x_k)$
 $N^{-2/3}$ -coupling: volume N , Coulomb interaction

$$\int_0^{N^{1/3}} x^{-1} d^3x \sim N^{2/3}$$

- ▶ Macroscopic equation: Hartree (Fock)
 $d_t \phi_t^j = \left(-\Delta + A + N^{-1} \sum_{k=1}^N |\phi_t^k|^2 \star V \right) \phi_t^j$
 $- \left(N^{-1} \sum_{k=1}^N \left(\phi_t^{k,*} \phi_t^j \right) \star V \right) \phi_t^k$
- ▶ $\rho = \sum_{j=1}^N |\phi_t^j\rangle \langle \phi_t^j|$
- ▶ Difficulties: $l \neq 0$
 $N^{-2/3}$ instead of N^{-1}

Fermions

- ▶ Microscopic system: $\Psi_0 = \Lambda_{j=1}^N \phi_0^j$
 $H = \sum_{j=1}^N -\Delta_j + \sum_{j=1}^N A_t(x_j) + N^{-2/3} \sum_{k < j} V(x_j - x_k)$
 $N^{-2/3}$ -coupling: volume N , Coulomb interaction

$$\int_0^{N^{1/3}} x^{-1} d^3x \sim N^{2/3}$$

- ▶ Macroscopic equation: Hartree (Fock)
 $d_t \phi_t^j = \left(-\Delta + A + N^{-1} \sum_{k=1}^N |\phi_t^k|^2 \star V \right) \phi_t^j$
 $- \left(N^{-1} \sum_{k=1}^N \left(\phi_t^{k,*} \phi_t^j \right) \star V \right) \phi_t^k$
- ▶ $\rho = \sum_{j=1}^N |\phi_t^j\rangle \langle \phi_t^j|$
- ▶ Difficulties: $I \neq 0$
 $N^{-2/3}$ instead of N^{-1}

Fermions

- ▶ Microscopic system: $\Psi_0 = \Lambda_{j=1}^N \phi_0^j$
 $H = \sum_{j=1}^N -\Delta_j + \sum_{j=1}^N A_t(x_j) + N^{-2/3} \sum_{k < j} V(x_j - x_k)$
 $N^{-2/3}$ -coupling: volume N , Coulomb interaction

$$\int_0^{N^{1/3}} x^{-1} d^3x \sim N^{2/3}$$

- ▶ Macroscopic equation: Hartree (Fock)
 $d_t \phi_t^j = \left(-\Delta + A + N^{-1} \sum_{k=1}^N |\phi_t^k|^2 \star V \right) \phi_t^j$
 $- \left(N^{-1} \sum_{k=1}^N \left(\phi_t^{k,*} \phi_t^j \right) \star V \right) \phi_t^k$
- ▶ $\rho = \sum_{j=1}^N |\phi_t^j\rangle \langle \phi_t^j|$
- ▶ Difficulties: $l \neq 0$
 $N^{-2/3}$ instead of N^{-1}

Results for Fermions

- ▶ Semiclassical situation: N. Benedikter, M. Porta, B. Schlein (2014)
- ▶ Quantum situation: S. Petrat, P.P. (2015); V. Bach, S. Breteaux, S. Petrat, P. P., T. Tzaneteas (2015)
- ▶ Problem: potential of leading order, force small!
- ▶ Goal: Consider coupling $N^{-1/3}$.

Results for Fermions

- ▶ Semiclassical situation: N. Benedikter, M. Porta, B. Schlein (2014)
- ▶ Quantum situation: S. Petrat, P.P. (2015); V. Bach, S. Breteaux, S. Petrat, P. P., T. Tzaneteas (2015)
- ▶ Problem: potential of leading order, force small!
- ▶ Goal: Consider coupling $N^{-1/3}$.

Results for Fermions

- ▶ Semiclassical situation: N. Benedikter, M. Porta, B. Schlein (2014)
- ▶ Quantum situation: S. Petrat, P.P. (2015); V. Bach, S. Breteaux, S. Petrat, P. P., T. Tzaneteas (2015)
- ▶ Problem: potential of leading order, force small!
- ▶ Goal: Consider coupling $N^{-1/3}$.

Results for Fermions

- ▶ Semiclassical situation: N. Benedikter, M. Porta, B. Schlein (2014)
- ▶ Quantum situation: S. Petrat, P.P. (2015); V. Bach, S. Breteaux, S. Petrat, P. P., T. Tzaneteas (2015)
- ▶ Problem: potential of leading order, force small!
- ▶ Goal: Consider coupling $N^{-1/3}$.

Results for Fermions

- ▶ Semiclassical situation: N. Benedikter, M. Porta, B. Schlein (2014)
- ▶ Quantum situation: S. Petrat, P.P. (2015); V. Bach, S. Breteaux, S. Petrat, P. P., T. Tzaneteas (2015)
- ▶ Problem: potential of leading order, force small!
- ▶ Goal: Consider coupling $N^{-1/3}$.

Special situation

- ▶ Tracer in filled Fermi sea: $\Psi_0 = \chi(y) \bigwedge_{j=1}^N \phi_j(x_j)$
- ▶ Interaction with tracer and gas particles: $H_I = \sum_{j=1}^N V(y, x_j)$
- ▶ Empirics: free evolution of tracer
- ▶ Strong contrast to bosonic case $\Psi_0 = \chi(y) \left(\prod_{j=1}^N \phi_j(x_j) \right)_{sym}$
Brownian motion.
- ▶ Mean field works much better in fermionic case!

Special situation

- ▶ Tracer in filled Fermi sea: $\Psi_0 = \chi(y) \bigwedge_{j=1}^N \phi_j(x_j)$
- ▶ Interaction with tracer and gas particles: $H_I = \sum_{j=1}^N V(y, x_j)$
- ▶ Empirics: free evolution of tracer
- ▶ Strong contrast to bosonic case $\Psi_0 = \chi(y) \left(\prod_{j=1}^N \phi_j(x_j) \right)_{sym}$
Brownian motion.
- ▶ Mean field works much better in fermionic case!

Special situation

- ▶ Tracer in filled Fermi sea: $\Psi_0 = \chi(y) \bigwedge_{j=1}^N \phi_j(x_j)$
- ▶ Interaction with tracer and gas particles: $H_I = \sum_{j=1}^N V(y, x_j)$
- ▶ Empirics: free evolution of tracer
- ▶ Strong contrast to bosonic case $\Psi_0 = \chi(y) \left(\prod_{j=1}^N \phi_j(x_j) \right)_{sym}$
Brownian motion.
- ▶ Mean field works much better in fermionic case!

Special situation

- ▶ Tracer in filled Fermi sea: $\Psi_0 = \chi(y) \bigwedge_{j=1}^N \phi_j(x_j)$
- ▶ Interaction with tracer and gas particles: $H_I = \sum_{j=1}^N V(y, x_j)$
- ▶ Empirics: free evolution of tracer
- ▶ Strong contrast to bosonic case $\Psi_0 = \chi(y) \left(\prod_{j=1}^N \phi_j(x_j) \right)_{sym}$
Brownian motion.
- ▶ Mean field works much better in fermionic case!

Special situation

- ▶ Tracer in filled Fermi sea: $\Psi_0 = \chi(y) \bigwedge_{j=1}^N \phi_j(x_j)$
- ▶ Interaction with tracer and gas particles: $H_I = \sum_{j=1}^N V(y, x_j)$
- ▶ Empirics: free evolution of tracer
- ▶ Strong contrast to bosonic case $\Psi_0 = \chi(y) \left(\prod_{j=1}^N \phi_j(x_j) \right)_{sym}$
Brownian motion.
- ▶ Mean field works much better in fermionic case!

Special situation

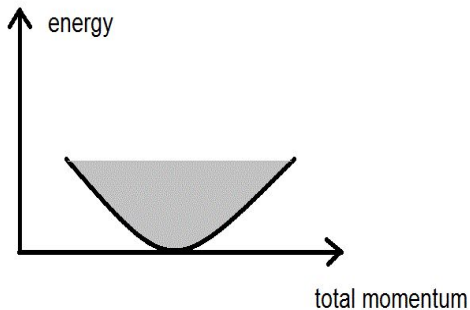
- ▶ Tracer in filled Fermi sea: $\Psi_0 = \chi(y) \bigwedge_{j=1}^N \phi_j(x_j)$
- ▶ Interaction with tracer and gas particles: $H_I = \sum_{j=1}^N V(y, x_j)$
- ▶ Empirics: free evolution of tracer
- ▶ Strong contrast to bosonic case $\Psi_0 = \chi(y) \left(\prod_{j=1}^N \phi_j(x_j) \right)_{sym}$
Brownian motion.
- ▶ Mean field works much better in fermionic case!

Special situation

- ▶ Tracer in filled Fermi sea: $\Psi_0 = \chi(y) \bigwedge_{j=1}^N \phi_j(x_j)$
- ▶ Interaction with tracer and gas particles: $H_I = \sum_{j=1}^N V(y, x_j)$
- ▶ Empirics: free evolution of tracer
- ▶ Strong contrast to bosonic case $\Psi_0 = \chi(y) \left(\prod_{j=1}^N \phi_j(x_j) \right)_{sym}$
Brownian motion.
- ▶ Mean field works much better in fermionic case!

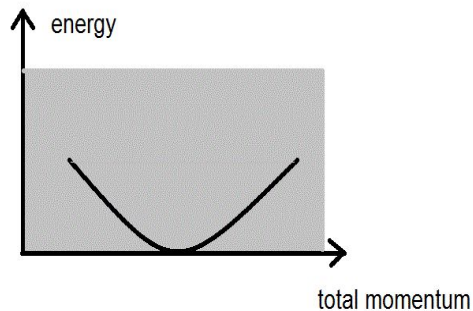
Special situation

1d: easy: Momentum and energy conservation.



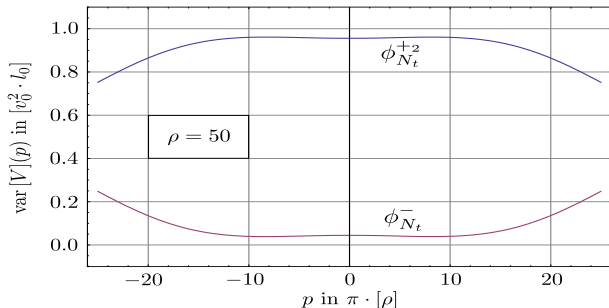
Special situation

Higher dimensions



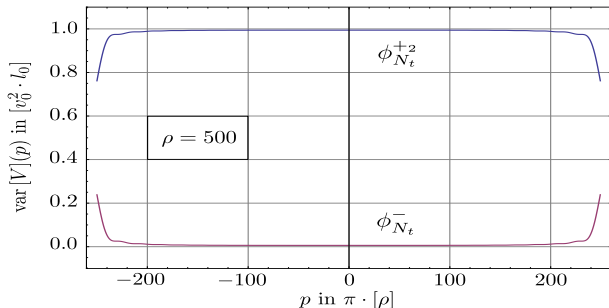
Estimate of fluctuations

Variance of force at some position y
(fermions: purple, bosons: blue)



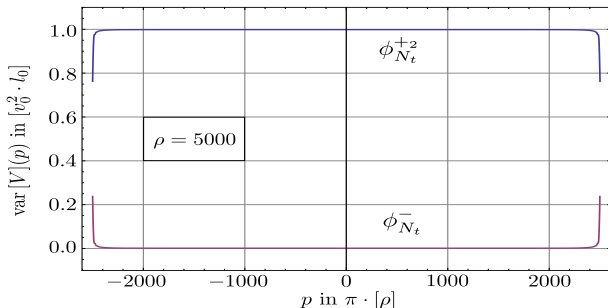
Estimate of fluctuations

Variance of force at some position y
(fermions: purple, bosons: blue)



Estimate of fluctuations

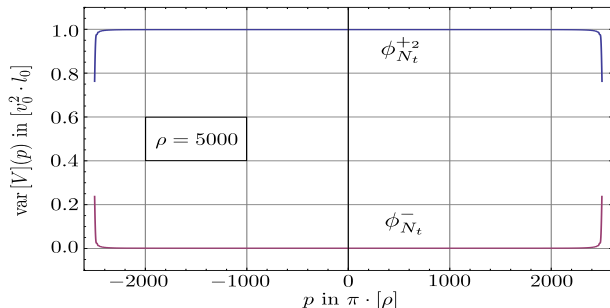
Variance of force at some position y
(fermions: purple, bosons: blue)



- ▶ Fluctuation of force is much smaller for fermions, still large
- ▶ Correlation due to antisymmetry reduces fluctuations.
- ▶ Fluctuations caused by particles with high momentum
Momentum transfer small.

Estimate of fluctuations

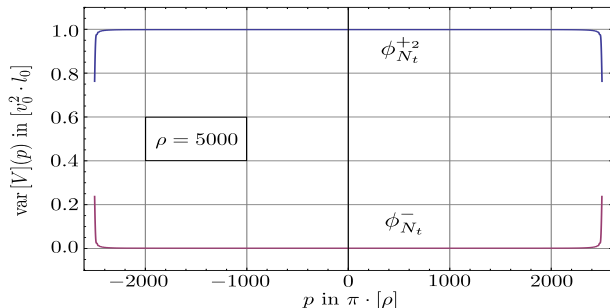
Variance of force at some position y
(fermions: purple, bosons: blue)



- ▶ Fluctuation of force is much smaller for fermions, still large
- ▶ Correlation due to antisymmetry reduces fluctuations.
- ▶ Fluctuations caused by particles with high momentum
Momentum transfer small.

Estimate of fluctuations

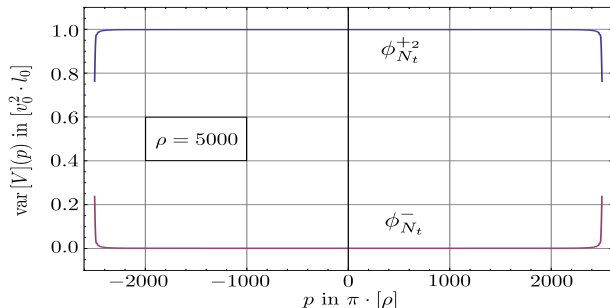
Variance of force at some position y
(fermions: purple, bosons: blue)



- ▶ Fluctuation of force is much smaller for fermions, still large
- ▶ Correlation due to antisymmetry reduces fluctuations.
- ▶ Fluctuations caused by particles with high momentum
Momentum transfer small.

Estimate of fluctuations

Variance of force at some position y
(fermions: purple, bosons: blue)

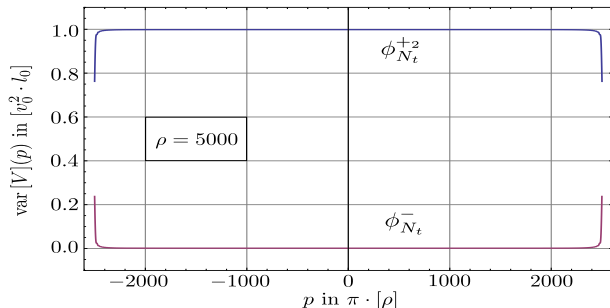


- ▶ Fluctuation of force is much smaller for fermions, still large
- ▶ Correlation due to antisymmetry reduces fluctuations.
- ▶ Fluctuations caused by particles with high momentum

Momentum transfer small.

Estimate of fluctuations

Variance of force at some position y
(fermions: purple, bosons: blue)



- ▶ Fluctuation of force is much smaller for fermions, still large
- ▶ Correlation due to antisymmetry reduces fluctuations.
- ▶ Fluctuations caused by particles with high momentum
Momentum transfer small.

Thank you!